

The spectral flow, the Fredholm index, and the spectral shift function

Alexander Pushnitski

Dedicated to M. Sh. Birman on the occasion of his 80th birthday

ABSTRACT. We discuss the well known “Fredholm index=spectral flow” theorem and show that it can be interpreted as a limit case of an identity involving two spectral shift functions.

1. Introduction

1.1. Background. Let $A(t)$, $t \in \mathbb{R}$, be a family of self-adjoint operators in a separable Hilbert space \mathcal{H} such that the limits

$$(1.1) \quad A^\pm = \lim_{t \rightarrow \pm\infty} A(t)$$

exist in an appropriate sense. In the Hilbert space $L^2(\mathbb{R}; \mathcal{H})$, consider the operator

$$(1.2) \quad D_A = \frac{d}{dt} + A(t), \text{ i.e. } (D_A u)(t) = \frac{du(t)}{dt} + A(t)u(t).$$

It is well known (see e.g. [4, 8] and references to earlier work therein) that, under the appropriate assumptions on $A(t)$, the Fredholm index of the operator D_A equals the spectral flow of the family $\{A(t)\}_{t \in \mathbb{R}}$ through zero. The spectral flow through zero should be understood as the number of eigenvalues of $A(t)$ (counting multiplicities) that cross zero from left to right minus the number of eigenvalues of $A(t)$ that cross zero from right to left as t grows from $-\infty$ to $+\infty$. The “Fredholm index = spectral flow through zero” theorem is one of the large family of index theorems; see e.g. [4] for discussion.

The “Fredholm index = spectral flow through zero” theorem is usually considered under the assumption that the spectra of the operators $A(t)$ are discrete, at least on some interval containing zero. The purpose of this note is to show that this assumption can be lifted at the expense of the trace class assumption

$$(1.3) \quad \int_{-\infty}^{\infty} \|A'(t)\|_{S_1} dt < \infty, \quad A'(t) \equiv \frac{dA(t)}{dt},$$

where $\|\cdot\|_{S_1}$ is the trace norm in \mathcal{H} , $\|A\|_{S_1} = \sqrt{\text{tr}(A^*A)}$. Assumption (1.3) ensures that the difference $A^+ - A^-$ is a trace class operator, which allows one to use the

notion of *M. G. Krein's spectral shift function* for the pair A^+, A^- . The point we would like to make is that the “Fredholm index = spectral flow” theorem can be understood as a particular limiting case of a fairly general identity (see (1.11) below) involving two spectral shift functions. This identity might be interesting in its own right.

It is a pleasure to dedicate this note to M. Sh. Birman, who has taught me (among many other useful things) to think of the spectral shift function whenever two self-adjoint operators are involved.

1.2. Notation. We denote by S_1 and S_2 the trace class and the Hilbert-Schmidt class, with the norms $\|\cdot\|_{S_1}$ and $\|\cdot\|_{S_2}$. For a self-adjoint operator A and an interval $\delta \subset \mathbb{R}$, we denote by $E_A(\delta)$ the spectral projection of A corresponding to δ . We denote by $N_A(\delta) = \text{Tr } E_A(\delta)$ the total number of eigenvalues (counting multiplicities) of A in the interval δ . For self-adjoint semi-bounded from below operators A and B , the inequality $A \leq B$ is understood in the quadratic form sense, i.e. for all sufficiently large $a > 0$, $\text{Dom}((A + a)^{1/2}) \supset \text{Dom}((B + a)^{1/2})$ and for all $f \in \text{Dom}((B + a)^{1/2})$, $\|(A + a)^{1/2}f\| \leq \|(B + a)^{1/2}f\|$.

1.3. The spectral shift function. Here we recall the necessary facts from the spectral shift function theory. See the original paper [5] or a survey [2] or a book [9] for the details.

Let H and \tilde{H} be self-adjoint operators in a Hilbert space. The simplest situation in which the spectral shift function can be defined is when the difference $\tilde{H} - H$ belongs to the trace class S_1 . Then there exists a unique function $\xi(\cdot; \tilde{H}, H) \in L^1(\mathbb{R})$ such that the Lifshits-Krein trace formula

$$(1.4) \quad \text{Tr}(f(\tilde{H}) - f(H)) = \int_{-\infty}^{\infty} \xi(\lambda; \tilde{H}, H) f'(\lambda) d\lambda$$

holds true for every $f \in C_0^\infty(\mathbb{R})$. The function $\xi(\cdot; \tilde{H}, H)$ is called the spectral shift function for the pair \tilde{H}, H . In fact, the class of admissible functions f in (1.4) is much wider than $C_0^\infty(\mathbb{R})$; see [2, 9] for the details and references to the literature.

The assumption $\tilde{H} - H \in S_1$ is very restrictive in applications. Suppose instead that \tilde{H} and H are non-negative (in the quadratic form sense) self-adjoint operators such that

$$(1.5) \quad (\tilde{H} - z)^{-1} - (H - z)^{-1} \in S_1$$

for some (and hence for all) $z \in \mathbb{C} \setminus [0, \infty)$. Then there exists a unique function $\xi(\cdot; \tilde{H}, H) \in L^1(\mathbb{R}, (1 + \lambda^2)^{-1} d\lambda)$, $\text{supp } \xi \in [0, \infty)$ such that the trace formula (1.4) holds true for all $f \in C_0^\infty(\mathbb{R})$.

Next, assuming either $\tilde{H} - H \in S_1$ or (1.5) holds true, suppose that for some (possibly semi-infinite) open interval $\Delta \subset \mathbb{R}$ we have $\sigma_{\text{ess}}(H) \cap \Delta = \emptyset$. By Weyl's theorem on the invariance of the essential spectrum with respect to compact perturbations, we also have $\sigma_{\text{ess}}(\tilde{H}) \cap \Delta = \emptyset$. Then, as it is not difficult to see from (1.4), for any $a, b \in \Delta \setminus (\sigma(H) \cup \sigma(\tilde{H}))$, $a < b$, we have

$$(1.6) \quad \xi(b; \tilde{H}, H) - \xi(a; \tilde{H}, H) = N_H(a, b) - N_{\tilde{H}}(a, b),$$

where $N_H(a, b)$ is the number of eigenvalues (counting multiplicities) of H in (a, b) . Formula (1.6) remains true in the case $a = -\infty$ (or $b = \infty$).

1.4. The spectral shift function and the spectral flow. Let $H(\alpha)$, $\alpha \in [0, 1]$, be a family of self-adjoint operators such that the operators $H(\alpha) - H(0)$ belong to the trace class and depend continuously on α in the trace norm. Then the function $\xi(\cdot; H(\alpha), H(0))$ is well defined and continuous in α as an element of $L^1(\mathbb{R})$.

As noted above, by Weyl's theorem $\sigma_{ess}(H(\alpha))$ is independent of $\alpha \in [0, 1]$. Suppose that $\sigma_{ess}(H(\alpha)) \cap \Delta = \emptyset$ for some interval $\Delta \subset \mathbb{R}$. Then we claim that for any $\lambda \in \Delta \setminus (\sigma(H(1)) \cup \sigma(H(0)))$, the spectral shift function $\xi(\lambda; H(1), H(0))$ equals the spectral flow of the family $H(\alpha)$ through λ as α grows from 0 to 1:

$$(1.7) \quad \begin{aligned} \xi(\lambda; H(1), H(0)) &= \langle \text{the number of eigenvalues of } H(\alpha) \text{ which cross } \lambda \text{ rightwards} \rangle \\ &\quad - \langle \text{the number of eigenvalues of } H(\alpha) \text{ which cross } \lambda \text{ leftwards} \rangle \end{aligned}$$

as long as the r.h.s. is finite.

In order to justify (1.7), first suppose that there exists $\lambda_0 < \lambda$ such that

$$(1.8) \quad \lambda_0 \in \Delta \setminus (\cup_{\alpha \in [0, 1]} \sigma(H(\alpha))).$$

Then it is not difficult to check that $\xi(\lambda_0; H(\alpha), H(0)) = 0$ for all α and so (1.6) (with $a = \lambda_0$, $b = \lambda$) yields

$$(1.9) \quad \xi(\lambda; H(\alpha), H(0)) = N_{H(0)}(\lambda_0, \lambda) - N_{H(\alpha)}(\lambda_0, \lambda).$$

Considering the r.h.s. of (1.9) as a function of α , it is easy to see that

$$\langle \text{r.h.s. of (1.9) with } \alpha = 1 \rangle = \langle \text{r.h.s. of (1.7)} \rangle$$

whenever the r.h.s. of (1.7) is well defined. Thus, (1.7) holds true.

In general, λ_0 as in (1.8) may not exist, but we can always split $[0, 1]$ into sufficiently small subintervals δ_i such that for each family $\{H(\alpha) \mid \alpha \in \delta_i\}$, λ_0 can be chosen appropriately. Then formula (1.7) can be obtained by adding up the formulas corresponding to all the subintervals.

1.5. Main result. It will be convenient to write $A(t) = A^- + B(t)$, where A^- is an arbitrary self-adjoint operator in \mathcal{H} and $B(t)$ is a family of trace class operators such that the derivative $B'(t) = \frac{dB(t)}{dt}$ exists in the trace norm and

$$(1.10) \quad \int_{-\infty}^{\infty} \|B'(t)\|_{S_1} dt < \infty.$$

This assumption ensures that the limits $B^\pm = \lim_{t \rightarrow \pm\infty} B(t)$ exist in trace norm. We assume that $B^- = 0$ (of course, this is merely a normalization condition) and define $A^+ = A^- + B^+$. According to this definition, for all $t \in \mathbb{R}$ the operators $A(t)$ have the same domain $\text{Dom}(A(t)) = \text{Dom}(A^-)$.

Consider the operator D_A (see (1.2)) in the Hilbert space $L^2(\mathbb{R}, \mathcal{H})$ with the domain $\text{Dom}(D_A)$ consisting of all u from the Sobolev space $W_2^1(\mathbb{R}, \mathcal{H})$ such that

$$u(t) \in \text{Dom}(A^-) \text{ for all } t \text{ and } \int_{-\infty}^{\infty} (\|u'(t)\|^2 + \|A^- u(t)\|^2) dt < \infty.$$

The operator D_A is closed. This can be seen as follows. Let us write D_A as a sum $D_A = D_{A^-} + B$. The operator B of multiplication by $B(t)$ is bounded in $L^2(\mathbb{R}; \mathcal{H})$, so the question reduces to the closedness of D_{A^-} . Note that $D_{A^-} = i(\frac{1}{i} \frac{d}{dt}) + A^-$, and the self-adjoint operators $(\frac{1}{i} \frac{d}{dt})$ and A^- in $L^2(\mathbb{R}; \mathcal{H})$ commute. Using the

spectral representations of $(\frac{1}{i}\frac{d}{dt})$ and A^- , it is easy to see that D_{A^-} is closed on the domain $\text{Dom}(D_{A^-}) = \text{Dom}(D_A)$.

The same argument shows that the adjoint operator D_A^* is defined as a closed operator on the same domain as D_A . Below we consider the self-adjoint operators

$$H = D_A^* D_A \quad \text{and} \quad \tilde{H} = D_A D_A^*.$$

As the difference $A^+ - A^-$ is a trace class operator, we can consider the spectral shift function $\xi(\lambda; A^+, A^-)$, $\lambda \in \mathbb{R}$. The first part of the Theorem below shows that the spectral shift function $\xi(\lambda; \tilde{H}, H)$ is also well defined. The following Theorem relates these two spectral shift functions.

THEOREM 1.1. *Assume (1.3) and let H, \tilde{H} be as defined above. For any $z \in \mathbb{C} \setminus [0, \infty)$, the difference $(\tilde{H} - z)^{-1} - (H - z)^{-1}$ belongs to the trace class. For a.e. $\lambda > 0$, we have the identity*

$$(1.11) \quad \xi(\lambda; \tilde{H}, H) = \frac{1}{\pi} \int_{-\sqrt{\lambda}}^{\sqrt{\lambda}} \xi(s; A^+, A^-) \frac{ds}{\sqrt{\lambda - s^2}},$$

where the integral in the r.h.s. converges absolutely.

COROLLARY 1.2. *Suppose that 0 is not in the spectrum of A^+ or A^- . Then the operator D_A is Fredholm and*

$$(1.12) \quad \text{index } D_A = \dim \text{Ker } D_A - \dim \text{Ker } D_A^* = \xi(0; A^+, A^-).$$

Generally speaking, the spectral shift function $\xi(\lambda; A^+, A^-)$ is defined as an element of $L^1(\mathbb{R})$, so it does not make sense to speak of its value at a fixed point $\lambda = 0$. However, the assumption of Corollary 1.2 implies that $\xi(\lambda; A^+, A^-)$ is constant near $\lambda = 0$, and so $\xi(0; A^+, A^-)$ is well defined.

According to (1.7), the r.h.s. of (1.12) coincides with the spectral flow of the family $A(t)$ through zero as long as the spectral flow is well defined. Thus, Corollary 1.2 can be interpreted as the generalisation of the “Fredholm index=spectral flow” theorem.

Corollary 1.2, under various sets of conditions on $A(t)$, is well known; see, e.g. [8, 4, 3] and references to earlier works therein.

1.6. The strategy of the proof. The proof of Theorem 1.1 is based on an identity due to [4, (3.14)] which we state below as Proposition 1.3. In order to state this identity, let us fix the principal branch of the square root in $\mathbb{C} \setminus (-\infty, 0]$. For any $z \in \mathbb{C} \setminus [0, \infty)$ and any $s \in \mathbb{R}$, we denote

$$g_z(s) = \frac{s}{\sqrt{s^2 - z}}.$$

The formula below involves traces of operators in \mathcal{H} and in $L^2(\mathbb{R}; \mathcal{H})$. We denote by tr the trace in \mathcal{H} and by Tr the trace in $L^2(\mathbb{R}; \mathcal{H})$.

PROPOSITION 1.3. *Assume (1.3) and let \tilde{H}, H be as defined above. Then for any $z \in \mathbb{C} \setminus [0, \infty)$, the difference $g_z(A^+) - g_z(A^-)$ belongs to the trace class in \mathcal{H} and*

$$(1.13) \quad \text{Tr}((\tilde{H} - z)^{-1} - (H - z)^{-1}) = \frac{1}{2z} \text{tr}(g_z(A^+) - g_z(A^-)).$$

The identity (1.13) was proven¹ in [4, (3.14)] for the case $\dim \mathcal{H} < \infty$ as a particular case of a more general trace identity. We will give a more streamlined proof of (1.13), based on the ideas from [3], where (1.13) was proven in the case $\dim \mathcal{H} = 1$. This plan of the proof is as follows. We first note that the operators H and \tilde{H} can be represented as

$$(1.14) \quad H = -\frac{d^2}{dt^2} + Q(t), \quad \tilde{H} = -\frac{d^2}{dt^2} + \tilde{Q}(t),$$

$$(1.15) \quad \text{where } Q(t) = A(t)^2 - A'(t), \quad \tilde{Q}(t) = A(t)^2 + A'(t).$$

Then, following [6], we express the integral kernels of the resolvents $(H - z)^{-1}$ and $(\tilde{H} - z)^{-1}$ in terms of the solutions to the operator differential equation $-F'' + QF = zF$. Integrating the traces of these resolvent kernels over the diagonal, we obtain the expression the l.h.s. of (1.13).

Given (1.13), one can derive Theorem 1.1 fairly easily by applying the Lifshits-Krein trace formula (1.4) to both sides of (1.4). This is done in Section 2. Proposition 1.3 in the case $\dim \mathcal{H} < \infty$ and Corollary 1.2 are also proven in Section 2.

In Section 3 (which is of a technical nature) we use an approximation argument to extend Proposition 1.3 from the case $\dim \mathcal{H} < \infty$ to the case $\dim \mathcal{H} = \infty$.

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2. Proofs

2.1. Trace class inclusions. Our first task is to show that the operators in both sides of (1.13) belong to the trace class.

LEMMA 2.1. *Under the assumptions of Theorem 1.1, the operators $(H - z)^{-1} - (\tilde{H} - z)^{-1}$ and $g_z(A^+) - g_z(A^-)$ belong to the trace class for any $z \in \mathbb{C} \setminus [0, \infty)$.*

Note that Lemma 2.1 in particular ensures that the spectral shift function $\xi(\cdot; \tilde{H}, H)$ is well defined.

Let us introduce some notation. We denote by H_0 the self-adjoint operator $-d^2/dt^2$ in $L^2(\mathbb{R}, \mathcal{H})$ and by $R_0(z) = (H_0 - z)^{-1}$ its resolvent. We also denote $R(z) = (H - z)^{-1}$ and $\tilde{R}(z) = (\tilde{H} - z)^{-1}$.

First, we need

LEMMA 2.2. *Let $V(t)$, $t \in \mathbb{R}$, be a family of trace class operators in \mathcal{H} such that*

$$(2.1) \quad \int_{-\infty}^{\infty} \|V(t)\|_{S_1} dt < \infty,$$

and let V be the operator in $L^2(\mathbb{R}, \mathcal{H})$, $(Vu)(t) = V(t)u(t)$. Then for any $z < 0$, one has

$$(2.2) \quad R_0(z)^{1/2} V R_0(z)^{1/2} \in S_1 \text{ and } \|R_0(z)^{1/2} V R_0(z)^{1/2}\|_{S_1} \leq \frac{1}{4\sqrt{|z|}} \int_{-\infty}^{\infty} \|V(t)\|_{S_1} dt.$$

¹This identity is stated in [4] with a wrong sign; compare with [3, Example 4.1]

In particular, V is a relatively form compact perturbation of H_0 .

PROOF. Write $V(t) = |V(t)|^{1/2} \text{sign}(V(t)) |V(t)|^{1/2}$. Using the Fourier transform, one easily checks that

$$|V|^{1/2} R_0(z)^{1/2} \in S_2$$

and

$$\| |V|^{1/2} R_0(z)^{1/2} \|_{S_2}^2 \leq \frac{1}{4\sqrt{|z|}} \int_{-\infty}^{\infty} \| |V(t)|^{1/2} \|_{S_2}^2 dt = \frac{1}{4\sqrt{|z|}} \int_{-\infty}^{\infty} \| V(t) \|_{S_1} dt.$$

It remains to write

$$\begin{aligned} \| R_0(z)^{1/2} V R_0(z)^{1/2} \|_{S_1} &= \| R_0(z)^{1/2} |V|^{1/2} \text{sign}(V) |V|^{1/2} R_0(z)^{1/2} \|_{S_1} \\ &\leq \| |V|^{1/2} R_0(z)^{1/2} \|_{S_2}^2, \end{aligned}$$

which completes the proof. \square

PROOF OF LEMMA 2.1.

1. First we consider the difference of resolvents $R(z) - \tilde{R}(z)$. By a well known argument, it suffices to prove that $R(z) - \tilde{R}(z) \in S_1$ for at least one value of z . Using Lemma 2.2, let us choose $z \in (-\infty, -1)$ with $|z|$ sufficiently large so that

$$(2.3) \quad \| R_0(z)^{1/2} B' R_0(z)^{1/2} \| \leq 1/2.$$

We have

$$\begin{aligned} H - z &= H_0 + A^2 - B' - z \geq H_0 - B' - z \\ &= (H_0 - z)^{1/2} (I - R_0(z)^{1/2} B' R_0(z)^{1/2}) (H_0 - z)^{1/2}; \end{aligned}$$

the inequality is understood in the sense of the quadratic forms. Thus we have

$$R(z) \leq R_0(z)^{1/2} (I - R_0(z)^{1/2} B' R_0(z)^{1/2})^{-1} R_0(z)^{1/2} \leq 2R_0(z).$$

It follows that the resolvent $R(z)$ can be represented as

$$(2.4) \quad R(z) = M(z) R_0(z)^{1/2} = R_0(z)^{1/2} M(z)^* \quad \text{with} \quad \| M(z) \|^2 \leq 2.$$

The same argument works for $\tilde{R}(z)$, so we have

$$(2.5) \quad \tilde{R}(z) = \tilde{M}(z) R_0^{1/2}(z) = R_0(z)^{1/2} \tilde{M}(z)^* \quad \text{with} \quad \| \tilde{M}(z) \|^2 \leq 2.$$

Now we can rewrite the resolvent identity as

$$R(z) - \tilde{R}(z) = 2R(z) B' \tilde{R}(z) = 2M(z) (R_0(z)^{1/2} B' R_0(z)^{1/2}) \tilde{M}(z)^*,$$

and, by Lemma 2.2, the r.h.s. belongs to the trace class.

2. Consider the difference $g_z(A^+) - g_z(A^-)$. We follow the well known argument (see e.g. [2]). Let \tilde{g}_z be the Fourier transform of the derivative of g_z ; then \tilde{g}_z is absolutely integrable and we can write

$$(2.6) \quad g_z(\lambda) = g_z(0) + i \int_{-\infty}^{\infty} \frac{e^{-i\lambda t} - 1}{t} \tilde{g}_z(t) dt \quad \text{with} \quad \int_{-\infty}^{\infty} |\tilde{g}_z(t)| dt < \infty.$$

Thus, we have the representation

$$(2.7) \quad g_z(A^+) - g_z(A^-) = \int_{-\infty}^{\infty} dt \tilde{g}_z(t) t^{-1} \int_0^s ds e^{-i(t-s)A^+} (A^+ - A^-) e^{-isA^-}.$$

Since

$$\| e^{-i(t-s)A^+} (A^+ - A^-) e^{-isA^-} \|_{S_1} \leq \| (A^+ - A^-) \|_{S_1},$$

the integral in the r.h.s of (2.7) converges in trace norm. \square

2.2. Proof of Theorem 1.1. 1. Denote for brevity $\xi(\lambda) = \xi(\lambda; \tilde{H}, H)$ and $\eta(\lambda) = \xi(\lambda; A^+, A^-)$. According to the Lifshits-Krein trace formula (1.4), the identity (1.13) can be rewritten as

$$-\int_0^\infty \frac{\xi(t)}{(t-z)^2} dt = \frac{1}{2z} \int_{-\infty}^\infty \eta(t) \left(\frac{\partial}{\partial t} g_z(t) \right) dt,$$

which can be rewritten as

$$\int_0^\infty \xi(t) \left(\frac{\partial}{\partial z} (t-z)^{-1} \right) dt = \int_{-\infty}^\infty \eta(t) \left(\frac{\partial}{\partial z} \frac{1}{\sqrt{t^2-z}} \right) dt.$$

Integrating over z , we get

$$(2.8) \quad \int_0^\infty \xi(t) \left(\frac{1}{t-z} - \frac{1}{t+1} \right) dt = \int_{-\infty}^\infty \eta(t) \left(\frac{1}{\sqrt{t^2-z}} - \frac{1}{\sqrt{t^2+1}} \right) dt.$$

2. Now we would like to take the imaginary parts of both sides of (2.8) and pass to the limit as $z \rightarrow \lambda + i0$ for a.e. $\lambda > 0$. By the well known properties of the Cauchy integrals, we have for a.e. $\lambda > 0$

$$\xi(\lambda) = \lim_{\varepsilon \rightarrow +0} \frac{1}{\pi} \int_0^\infty \xi(t) \operatorname{Im} \frac{1}{t - \lambda - i\varepsilon} dt.$$

Now consider the r.h.s. of (2.8). As $\eta \in L^1(\mathbb{R})$, it is easy to see that the integral

$$\int_{-\infty}^\infty \frac{|\eta(t)|}{|t^2 - \lambda|^{1/2}} dt$$

converges for a.e. $\lambda > 0$. By the dominated convergence theorem, this ensures that

$$\operatorname{Im} \int_{-\infty}^\infty \frac{\eta(t)}{\sqrt{t^2 - \lambda - i\varepsilon}} dt \rightarrow \int_{-\sqrt{\lambda}}^{\sqrt{\lambda}} \frac{\eta(\lambda)}{\sqrt{\lambda - t^2}} dt, \quad \varepsilon \rightarrow +0$$

for a.e. $\lambda > 0$. This allows us to pass to the limit in (2.8), which yields the required result. \square

2.3. Proof of Corollary 1.2. 1. In order to check that D_A is Fredholm, we use the following necessary and sufficient condition from [1, Theorem A.4]: *A closed operator T is Fredholm if and only if $0 \notin \sigma_{ess}(T^*T)$ and $0 \notin \sigma_{ess}(TT^*)$ and then $\operatorname{index}(T) = \dim \operatorname{Ker}(T^*T) - \dim \operatorname{Ker}(TT^*)$.*

Using our assumptions (1.3) and $0 \in \rho(A^-) \cap \rho(A^+)$, we can find $a > 0$ such that for all sufficiently large $|t|$, one has $A(t)^2 \geq aI$. One has

$$H = H_0 + A(t)^2 - B'(t) \geq H_0 + aI + V(t),$$

where

$$V(t) = -B'(t) + (A(t)^2 - aI)E_{A(t)^2}([0, a]).$$

The operator $(A(t)^2 - aI)E_{A(t)^2}([0, a])$ is of a finite rank for all $t \in \mathbb{R}$ and vanishes for all sufficiently large $|t|$. Thus, V satisfies (2.1) and therefore, by Lemma 2.2, V is a relatively form compact perturbation of H_0 . It follows that $\sigma_{ess}(H_0 + aI + V) = \sigma_{ess}(H_0 + aI) = [a, \infty)$. Thus, $\inf \sigma_{ess}(H) \geq a > 0$. The same argument applies to \tilde{H} . By the necessary and sufficient condition quoted above, D_A is a Fredholm operator and

$$(2.9) \quad \operatorname{index} D_A = \dim \operatorname{Ker} H - \dim \operatorname{Ker} \tilde{H}.$$

2. By the previous step, the spectra of H and \tilde{H} are discrete on $[0, a)$. It is well known (see e.g. [1, Lemma A.3]) that, since $H = D_A^* D_A$ and $\tilde{H} = D_A D_A^*$, one has $\sigma(H) \setminus \{0\} = \sigma(\tilde{H}) \setminus \{0\}$ and the multiplicities of the eigenvalues coincide. Thus, from (1.6) we get

$$(2.10) \quad \xi(\lambda; \tilde{H}, H) = \dim \operatorname{Ker} H - \dim \operatorname{Ker} \tilde{H}, \quad \lambda \in (0, a).$$

This identity and its relation to index of D_A is due to [3].

3. Combining Theorem 1.1 with (2.9) and (2.10), we obtain

$$\operatorname{index} D_A = \frac{1}{\pi} \int_{-\sqrt{\lambda}}^{\sqrt{\lambda}} \frac{\xi(t; A^+, A^-)}{\sqrt{\lambda - t^2}} dt, \quad \lambda \in (0, a).$$

Since $0 \in \rho(A^-) \cap \rho(A^+)$, the function $\xi(t; A^+, A^-)$ is constant near $t = 0$. Taking into account the identity

$$\frac{1}{\pi} \int_{-\sqrt{\lambda}}^{\sqrt{\lambda}} \frac{1}{\sqrt{\lambda - t^2}} dt = 1,$$

we arrive at (1.12). \square

2.4. Proof of Proposition 1.3 for “nice” $A(t)$. Our proof of Proposition 1.3 consists of two steps: we first prove the identity (1.13) for very “nice” functions $A(t)$ and then use approximation argument. Here we present the first step; the approximation argument, which has a more technical nature, is given in section 3.

LEMMA 2.3. *Suppose that $\dim \mathcal{H} < \infty$. Assume also that $A(t) = A^\pm$ for all sufficiently large $\pm t > 0$. Then Proposition 1.3 holds true.*

PROOF. 1. Of course, in the finite dimensional case all operators belong to the trace class, so we only need to prove the identity (1.13). In the case $\dim \mathcal{H} = 1$, this formula has been proven in [3, Example 4.1]. Below is a direct generalisation of the argument of [3] to the matrix-valued case. We call elements of \mathcal{H} “matrices”.

By analyticity in z it suffices to consider the case of real negative z . In what follows, we assume $z \in (-\infty, 0)$ and denote

$$\kappa_\pm = \sqrt{(A^\pm)^2 - z}.$$

The matrices κ_\pm in \mathcal{H} are positive definite.

2. We will compute the trace in the l.h.s. of (1.13) by constructing the integral kernels of the resolvents of H and \tilde{H} and evaluating integrals of these kernels over the diagonal. Our first aim is to construct the resolvent of H in terms of the solutions to the matrix Schrodinger equation. Here we follow [6].

Consider the matrix valued solutions $F_\pm(t)$, $t \in \mathbb{R}$, to the equation (see (1.14), (1.15))

$$(2.11) \quad -F_\pm'' + QF_\pm = zF_\pm,$$

satisfying the asymptotic conditions

$$(2.12) \quad F_\pm(t) = e^{\mp \kappa_\pm t}, \quad \pm t > 0 \text{ large}.$$

Existence of solutions F_\pm can be proven in the usual way by converting (2.11), (2.12) into Volterra type integral equations.

We have the relations

$$(2.13) \quad \begin{aligned} F_+(t) &= e^{-\varkappa_- t} a + e^{\varkappa_- t} b, \quad -t > 0 \text{ large}, \\ F_-(t) &= e^{\varkappa_+ t} c + e^{-\varkappa_+ t} d, \quad t > 0 \text{ large}, \end{aligned}$$

for some matrices a, b, c, d (which, of course, depend on $z < 0$). From here it is straightforward to see that if either a or c has a non-trivial kernel, then z is an eigenvalue of H . But, by definition, H cannot have negative eigenvalues, so we have

$$(2.14) \quad \text{Ker } a = \text{Ker } c = \{0\}$$

for all $z < 0$.

3. For any two solutions F, G to the equation (2.11), let us define the Wronskian

$$W(F, G) = F(t)^* G'(t) - F'(t)^* G(t).$$

By a direct calculation, the Wronskian does not depend on t . In particular, using the limiting forms (2.13) of the solutions F_{\pm} , we obtain

$$(2.15) \quad W(F_+, F_-) = 2\varkappa_+ c = 2a^* \varkappa_-.$$

By (2.14), it follows that $W(F_+, F_-)$ is invertible (it is here that we need $\dim \mathcal{H} < \infty$). Now we can construct the integral kernel $R(t, s)$ of the resolvent $(H - z)^{-1}$ as in [6]. We have

$$(2.16) \quad R(t, s) = \begin{cases} F_+(t)(W(F_+, F_-)^*)^{-1} F_-^*(s), & t \geq s, \\ F_-(t)(W(F_+, F_-))^{-1} F_+^*(s), & t < s. \end{cases}$$

4. As above, we can construct the integral kernel $\tilde{R}(x, y)$ of the resolvent $(\tilde{H} - z)^{-1}$ in terms of the solutions $\tilde{F}_{\pm}(t)$ to

$$(2.17) \quad -\tilde{F}_{\pm}'' + \tilde{Q}\tilde{F}_{\pm} = z\tilde{F}_{\pm}, \text{ and } \tilde{F}_{\pm}(t) = e^{\mp \varkappa_{\pm} t}, \quad \pm t > 0 \text{ large}.$$

By a direct calculation, the function

$$\tilde{F}_{\pm}(t) = (F'_{\pm}(t) + A(t)F_{\pm}(t))(\mp \varkappa_{\pm} + A^{\pm})^{-1}$$

satisfies (2.17). Also by a direct calculation,

$$W(\tilde{F}_+, \tilde{F}_-) = z(-\varkappa_+ + A^+)^{-1} W(F_+, F_-)(\varkappa_- + A^-)^{-1}.$$

This allows us to compute the kernel $\tilde{R}(t, s)$ in terms of the solutions F_{\pm} :

$$(2.18) \quad R(t, s) = \begin{cases} \frac{1}{z}(F'_+(t) + A(t)F_+(t))(W(F_+, F_-)^*)^{-1}(F'_-(s) + A(s)F_-(s))^*, & t \geq s, \\ \frac{1}{z}(F'_-(t) + A(t)F_-(t))(W(F_+, F_-))^{-1}(F'_+(s) + A(s)F_+(s))^*, & t < s. \end{cases}$$

5. Now we are ready to compute the trace in the l.h.s. of (1.13):

$$\text{Tr}((\tilde{H} - z)^{-1} - (H - z)^{-1}) = \lim_{R \rightarrow \infty} \int_{-R}^R \text{tr}(\tilde{R}(t, t) - R(t, t)) dt.$$

Using our formulas (2.16), (2.18) for the kernels $R(t, t)$ and $\tilde{R}(t, t)$, we obtain

$$\text{tr}(\tilde{R}(t, t) - R(t, t)) = \text{tr}(W(F_+, F_-)^*)^{-1} \left(\frac{1}{z} ((F'_-)^* + F_-^* A) (F'_+ + A F_+) - F_-^* F_+ \right).$$

Integrating by parts, after a little algebra we get

$$(2.19) \quad \int_{-R}^R \text{tr}(\tilde{R}(t, t) - R(t, t)) dt = \frac{1}{z} \text{tr}((W(F_+, F_-)^*)^{-1} F_-^* (F'_+ + A F_+)) \Big|_{-R}^R.$$

Now we can calculate the r.h.s. of (2.19) for large $R > 0$, using formula (2.15) and the asymptotic forms (2.12), (2.13) of F_{\pm} :

$$\begin{aligned} \operatorname{tr}(W(F_+, F_-)^{-1} F_-^* F'_+)|_{-R}^R &= -\frac{1}{2} \operatorname{tr}((c^*)^{-1} d^* e^{-2\kappa_+ R}) - \frac{1}{2} \operatorname{tr}(ba^{-1} e^{-2\kappa_- R}), \\ \operatorname{tr}(W(F_+, F_-)^{-1} F_-^* A F_+)|_{-R}^R &= \frac{1}{2} \operatorname{tr}(\kappa_+^{-1} A^+ - \kappa_-^{-1} A^-) \\ &\quad + \frac{1}{2} \operatorname{tr}(\kappa_+^{-1} (c^*)^{-1} d^* A^+ e^{-2\kappa_+ R}) - \frac{1}{2} \operatorname{tr}(ba^{-1} \kappa_-^{-1} A^- e^{-2\kappa_- R}). \end{aligned}$$

As κ_{\pm} are positive definite matrices, we have $\|e^{-2\kappa_{\pm} R}\| \rightarrow 0$ as $R \rightarrow \infty$. Thus,

$$\operatorname{Tr}((\tilde{H} - z)^{-1} - (H - z)^{-1}) = \frac{1}{2z} \operatorname{tr}(\kappa_+^{-1} A^+ - \kappa_-^{-1} A^-),$$

as required. \square

3. Approximation argument

First we give a general statement about convergence in both sides of the identity (1.13) and then construct the approximating sequence $A_n(t)$.

3.1. Convergence in (1.13). Let $A(t) = A^- + B(t)$ and $A_n(t) = A_n^- + B_n(t)$ be operator families satisfying (1.10). As above, we assume $B(-\infty) = B_n(-\infty) = 0$ and define $A_n^+ = A_n^- + B_n(+\infty)$. We also define $H_n = D_{A_n}^* D_{A_n}$, $\tilde{H}_n = D_{A_n} D_{A_n}^*$, and

$$R_n(z) = (H_n - z)^{-1}, \quad \tilde{R}_n(z) = (\tilde{H}_n - z)^{-1}.$$

LEMMA 3.1. *Assume that $\operatorname{Dom} A^- \subset \operatorname{Dom} A_n^-$ for all n and $A_n^- f \rightarrow A^- f$ for all $f \in \operatorname{Dom}(A^-)$. Next, assume that*

$$(3.1) \quad \int_{-\infty}^{\infty} \|B'_n(t) - B'(t)\|_{S_1} dt \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Then, for all $z \in (-\infty, -1)$ with sufficiently large $|z|$, one has

$$(3.2) \quad \|(g_z(A^+) - g_z(A^-)) - (g_z(A_n^+) - g_z(A_n^-))\|_{S_1} \rightarrow 0,$$

$$(3.3) \quad \|(R(z) - \tilde{R}(z)) - (R_n(z) - \tilde{R}_n(z))\|_{S_1} \rightarrow 0.$$

REMARK 3.2. 1. Assumption (3.1) implies that

$$(3.4) \quad \|(A_n^+ - A_n^-) - (A^+ - A^-)\|_{S_1} \rightarrow 0.$$

2. Our assumptions on A^- , A_n^- imply that $A_n^- \rightarrow A^-$ in strong resolvent sense; see [7, Theorem VIII.25(a)] and its proof. Combining this with (3.4), we see that also $A_n^+ \rightarrow A^+$ in strong resolvent sense.

We will repeatedly make use of the following well known fact, which holds true for any Schatten-von Neumann class S_p , $p \geq 1$, although we will only need it for the case $p = 1, 2$:

PROPOSITION 3.3. *Let T_n be a sequence of bounded operators in a Hilbert space which converges strongly to zero and let $M \in S_p$; then $\|T_n M\|_{S_p} \rightarrow 0$. If T_n^* also converges strongly to zero, then $\|MT_n\|_{S_p} \rightarrow 0$.*

The first part of this Proposition can be found, for example, in [9, Lemma 6.1.3], and the second part follows immediately by conjugation, since $\|MT_n\|_{S_p} = \|T_n^* M^*\|_{S_p}$.

PROOF OF (3.2). Writing the representation (2.7), we get

$$(3.5) \quad (g_z(A^+) - g_z(A^-)) - (g_z(A_n^+) - g_z(A_n^-)) = \int_{-\infty}^{\infty} dt \tilde{g}_z(t) t^{-1} \int_0^t ds K_n(t, s),$$

$$K_n(t, s) = e^{-i(t-s)A^+} (A^+ - A^-) e^{-isA^-} - e^{-i(t-s)A_n^+} (A_n^+ - A_n^-) e^{-isA_n^-}.$$

We would like to use the dominated convergence theorem in order to prove that the r.h.s. of (3.5) converges to zero in the trace norm. First note that

$$\|K_n(t, s)\|_{S_1} \leq \|A^+ - A^-\|_{S_1} + \|A_n^+ - A_n^-\|_{S_1}$$

and, by (3.4), the r.h.s. is bounded uniformly in n by some constant C . This gives an integrable bound for the integrand in the r.h.s. of (3.5).

Next, we claim that $\|K_n(t, s)\|_{S_1} \rightarrow 0$ for all t, s . Indeed, we can write

$$(3.6) \quad K_n(t, s) = (e^{-i(t-s)A^+} - e^{-i(t-s)A_n^+})(A^+ - A^-)e^{-isA^-} \\ + e^{-i(t-s)A_n^+}(A^+ - A^-)(e^{-isA^-} - e^{-isA_n^-}) + e^{-i(t-s)A_n^+}(A^+ - A^- - A_n^+ + A_n^-)e^{-isA_n^-}.$$

The last term in the r.h.s. converges to zero by (3.4). Next, since $A_n^\pm \rightarrow A^\pm$ in strong resolvent sense (see Remark 3.2), by [7, Theorem VIII.21] we have strong convergence $e^{itA_n^\pm} \rightarrow e^{itA^\pm}$ as $n \rightarrow \infty$. Thus, by Proposition 3.3, the first two terms in the r.h.s. of (3.6) converge to zero in the trace norm. By dominated convergence, this proves (3.2). \square

The proof of (3.3) requires a little more work. First we need an abstract lemma which ensures the strong convergence of resolvents. I am indebted to Nikolai Filonov for providing the proof of this lemma.

Let D be a closed densely defined operator in a Hilbert space such that D^* is also densely defined. For each n , let D_n be a closed densely defined operator such that D_n^* is also densely defined and $\text{Dom } D \subset \text{Dom } D_n$ and $\text{Dom } D^* \subset \text{Dom } D_n^*$.

LEMMA 3.4. *Assume the above conditions and assume that for all $f \in \text{Dom } D$ one has $\|D_n f - D f\| \rightarrow 0$ and for all $f \in \text{Dom } D^*$ one has $\|D_n^* f - D^* f\| \rightarrow 0$ as $n \rightarrow \infty$. Then for all $z \in \mathbb{C} \setminus [0, \infty)$, one has the strong convergence*

$$(3.7) \quad (D_n^* D_n - z)^{-1} \rightarrow (D^* D - z)^{-1}, \quad n \rightarrow \infty.$$

PROOF. 1. By [7, Chapter VIII, Problem 20], it suffices to prove the weak convergence in (3.7).

2. Fix $z \in \mathbb{C} \setminus [0, \infty)$, $f \in \mathcal{H}$ and denote $\varphi_n = (D_n^* D_n - z)^{-1} f$. We need to prove that $\varphi_n \rightarrow \varphi$ weakly, where $\varphi = (D^* D - z)^{-1} f$. First note that

$$(3.8) \quad \|\varphi_n\| \leq \frac{\|f\|}{\text{dist}(z, [0, \infty))}.$$

Next, we have

$$\|D_n \varphi_n\|^2 - z \|\varphi_n\|^2 = (f, \varphi_n) = (f, (D_n^* D_n - z)^{-1} f),$$

and so

$$(3.9) \quad \|D_n \varphi_n\|^2 \leq |z| \|\varphi_n\|^2 + |(f, (D_n^* D_n - z)^{-1} f)| \leq C(z) \|f\|^2.$$

By (3.8), (3.9) and the weak compactness of the unit ball in a Hilbert space, from the sequence φ_n one can choose a subsequence φ_{n_k} such that $\varphi_{n_k} \rightarrow \tilde{\varphi}$ and $D_{n_k} \varphi_{n_k} \rightarrow \psi$ weakly for some elements $\tilde{\varphi}, \psi$ in \mathcal{H} .

3. Let us prove that $\tilde{\varphi} \in \text{Dom } D$ and $\psi = D\tilde{\varphi}$. For any $\chi \in \text{Dom } D^*$, we have

$$(D_{n_k}\varphi_{n_k}, \chi) \rightarrow (\psi, \chi),$$

and so

$$(\varphi_{n_k}, D_{n_k}^*\chi) \rightarrow (\psi, \chi).$$

Since $\|D_{n_k}^*\chi - D^*\chi\| \rightarrow 0$ by our assumptions, we get

$$(3.10) \quad (\tilde{\varphi}, D^*\chi) = (\psi, \chi)$$

for all $\chi \in \text{Dom } D^*$; it follows that $\tilde{\varphi} \in \text{Dom } D$ and $\psi = D\tilde{\varphi}$.

4. Next, we have for any $\chi \in \text{Dom } D$:

$$(3.11) \quad (D_n\varphi_n, D_n\chi) - z(\varphi_n, \chi) = (f, \chi).$$

By the previous step, $\varphi_{n_k} \rightarrow \tilde{\varphi}$ and $D_{n_k}\varphi_{n_k} \rightarrow \psi$ weakly, and by the hypothesis, $D_n\chi \rightarrow D\chi$ strongly. Passing to the limit in (3.11) over the subsequence n_k , we get

$$(\psi, D\chi) - z(\tilde{\varphi}, \chi) = (f, \chi)$$

for all $\chi \in \text{Dom } D$. It follows that $\psi \in \text{Dom } D^*$ and $D^*\psi - z\tilde{\varphi} = f$. Recalling that $\psi = D\tilde{\varphi}$, we get that $\tilde{\varphi} \in \text{Dom}(D^*D)$ and $(D^*D - z)\tilde{\varphi} = f$. Thus, we have proven that $\tilde{\varphi} = \varphi$.

5. We have proven weak convergence $\varphi_n \rightarrow \varphi$ over a subsequence n_k . But we could have started from an arbitrary subsequence of φ_n and proven that it has a subsubsequence which weakly converges to φ . This proves that actually the whole sequence φ_n weakly converges to φ . \square

PROOF OF (3.3). 1. By Lemma 2.2 and the uniform boundedness of the integrals $\int \|B'_n(t)\|_{S_1} dt$, one can choose $a < -1$ such that for all $z \in (-\infty, a)$, the estimates

$$\|R_0(z)^{1/2}B'R_0(z)^{1/2}\| \leq 1/2, \quad \sup_n \|R_0(z)^{1/2}B'_nR_0(z)^{1/2}\| \leq 1/2$$

hold true. Then, as in the proof of Lemma 2.1, we get (2.4), (2.5), and also

$$\begin{aligned} R_n(z) &= M_n(z)R_0(z)^{1/2} = R_0(z)^{1/2}M_n(z)^*, \\ \tilde{R}_n(z) &= \tilde{M}_n(z)R_0(z)^{1/2} = R_0(z)^{1/2}\tilde{M}_n(z)^* \end{aligned}$$

with $\|M_n(z)\|^2 \leq 2$ and $\|\tilde{M}_n(z)\|^2 \leq 2$. In what follows, we fix $z \in (-\infty, a)$ as above and suppress the dependance of z in our notation for brevity.

2. Note that by Lemma 3.4, we have the strong convergence of resolvents $R_n \rightarrow R$, $\tilde{R}_n \rightarrow \tilde{R}$. Moreover, we claim that for the operators M_n, \tilde{M}_n we have the strong convergence $M_n \rightarrow M$, $\tilde{M}_n \rightarrow \tilde{M}$. Indeed, for all $f \in \text{Dom}(H_0 - z)^{1/2}$ one has

$$M_nf = R_n(H_0 - z)^{1/2}f \rightarrow R(H_0 - z)^{1/2}f = Mf.$$

Since the norms of M_n are uniformly bounded, we get the strong convergence $M_n \rightarrow M$. In the same way, we obtain the strong convergence $\tilde{M}_n \rightarrow \tilde{M}$.

3. Using the resolvent identity, we obtain

$$(3.12) \quad (R_n - \tilde{R}_n) - (R - \tilde{R}) = R_nB'_n\tilde{R}_n - RB'\tilde{R} \\ = R_n(B'_n - B')\tilde{R}_n + (R_n - R)B'\tilde{R}_n + RB'(\tilde{R}_n - \tilde{R}).$$

Let us consider separately each of the three terms in the r.h.s. of (3.12).

4. For the first term, we have

$$R_n(B'_n - B')\tilde{R}_n = M_n(R_0^{1/2}(B'_n - B')R_0^{1/2})\tilde{M}_n^*$$

and the r.h.s. converges to zero in the trace norm by Lemma 2.2 and assumption (3.1).

5. Consider the second term in the r.h.s. of (3.12). We have:

$$(R_n - R)B'\tilde{R}_n = (M_n - M)(R_0^{1/2}B'R_0^{1/2})\tilde{M}_n^*.$$

Since $R_0^{1/2}B'R_0^{1/2}$ is a trace class operator, and $M_n \rightarrow M$ strongly as $n \rightarrow \infty$, by Proposition 3.3, we obtain that the r.h.s. converges to zero in the trace norm.

6. Finally, the third term in the r.h.s. of (3.12) can be considered similarly to the second one:

$$[RB'(\tilde{R}_n - \tilde{R})]^* = (\tilde{R}_n - \tilde{R})B'R = (\tilde{M}_n - \tilde{M})(R_0^{1/2}B'R_0^{1/2})M^*,$$

and the r.h.s. goes to zero in the trace norm as $n \rightarrow \infty$. \square

3.2. Constructing the approximating family $A_n(t)$. We will approximate $A(t)$ in two steps. First, we approximate an arbitrary finite rank family $A(t)$ by the ones with compactly supported $A'(t)$. Next, we approximate an arbitrary family by finite rank families.

LEMMA 3.5. *Let $\dim \mathcal{H} < \infty$. Then Proposition 1.3 holds true.*

PROOF. By analyticity in z , it suffices to prove (1.13) for $z \in (-\infty, -1)$ with sufficiently large $|z|$. Then we can use Lemma 3.1.

For a given family $A(t)$, let us construct a sequence of families $A_n(t)$ such that for each n and all sufficiently large $\pm t > 0$, we have $A_n(t) = A^\pm$. This is not difficult to do. Indeed, let $A_n(t)$ be such that $A_n(t) = A(t)$ for $t \in [-n, n]$, $A_n(t) = A^-$ for $t \leq -n-1$, $A_n(t) = A^+$ for $t \geq n+1$, and $A_n(t)$ is obtained by linear interpolation on $[-n-1, -n]$ and $[n, n+1]$. Explicitly,

$$\begin{aligned} A_n(t) &= A^- + (t+n+1)(A(-n) - A^-), \quad t \in [-n-1, -n], \\ A_n(t) &= (t-n)A^+ + (n+1-t)A(n), \quad t \in [n, n+1]. \end{aligned}$$

Then we have

$$\begin{aligned} \int_{-\infty}^{\infty} \|A'_n(t) - A'(t)\|_{S_1} dt &\leq \int_n^{\infty} (\|A'_n(t)\|_{S_1} + \|A'(t)\|_{S_1}) dt \\ &\quad + \int_{-\infty}^{-n} (\|A'_n(t)\|_{S_1} + \|A'(t)\|_{S_1}) dt \rightarrow 0 \end{aligned}$$

as $n \rightarrow \infty$. By Lemma 2.3, the identity (1.13) holds true for the families $A_n(t)$. By Lemma 3.1, we can pass to the limit as $n \rightarrow \infty$ in both sides of (1.13), which yields the required result. \square

Next, we approximate an arbitrary family $A(t)$ by finite rank families.

LEMMA 3.6. *There exists a sequence of finite rank orthogonal projections P_n in \mathcal{H} such that:*

- (i) $P_n \rightarrow I$ strongly as $n \rightarrow \infty$;
- (ii) $\text{Ran } P_n \subset \text{Dom } A^-$ for all n ;
- (iii) for all $f \in \text{Dom}(A^-)$, one has $\|P_n A^- P_n f - A^- f\| \rightarrow 0$ as $n \rightarrow \infty$.

PROOF. For any $k \in \mathbb{Z}$, let E_k be the spectral projection of the operator A^- associated with the interval $(k, k+1]$, and let $Q_j^{(k)}$ be a sequence of finite rank projections such that $\text{Ran } Q_j^{(k)} \subset \text{Ran } E_k$ and $Q_j^{(k)} \rightarrow E_k$ strongly as $j \rightarrow \infty$. Take $P_n = \sum_{k=-n}^n Q_n^{(k)}$. Then (i), (ii) are obvious. Let us prove (iii). If $f \in \text{Dom } A^-$, then $f = \sum_{k=-\infty}^{\infty} f_k$, where $f_k = E_k f$ and $\sum_{k=-\infty}^{\infty} (k^2 + 1) \|f_k\|^2 < \infty$. Given such f and $\varepsilon > 0$, we can choose N sufficiently large so that $\sum_{|k| \geq N} (k^2 + 1) \|f_k\|^2 < \varepsilon^2$; denote $g_1 = \sum_{|k| < N} f_k$ and $g_2 = \sum_{|k| \geq N} f_k$. Then $\|P_n A^- P_n g_2\| \leq \varepsilon$ and $\|A^- g_2\| \leq \varepsilon$. On the other hand, it is easy to see that $\|P_n A^- P_n g_1 - A^- g_1\| \rightarrow 0$ as $n \rightarrow \infty$. This proves (iii). \square

PROOF OF PROPOSITION 1.3. 1. Let $A_n(t) = P_n A(t) P_n$, $B_n(t) = P_n B(t) P_n$, where P_n are as constructed in Lemma 3.6. We claim that

$$\int_{-\infty}^{\infty} \|B'_n(t) - B'(t)\|_{S_1} dt \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Indeed, let us apply the dominated convergence theorem. First, we have

$$\|B'_n(t) - B'(t)\|_{S_1} \leq \|B'_n(t)\|_{S_1} + \|B'(t)\|_{S_1} \leq 2\|B'(t)\|_{S_1}.$$

Next, for all $t \in \mathbb{R}$, we have

$$B'_n(t) - B'(t) = (P_n - I)B'(t)P_n + B'(t)(P_n - I),$$

and the right hand side converges to zero by Proposition 3.3, since $P_n \rightarrow I$ and $B'(t) \in S_1$.

2. By Lemma 3.5, the identity (1.13) holds true for the families $A_n(t)$. By Lemma 3.1, we can pass to the limit as $n \rightarrow \infty$ in both sides of (1.13) when $z \in (-\infty, -1)$, $|z|$ large. By analyticity, this yields the required result for all z . \square

References

- [1] J. Avron, R. Seiler and B. Simon, *The index of a pair of projections*, J. Funct. Anal. **120** (1994), no. 1, 220–237.
- [2] M. Sh. Birman and D. R. Yafaev, *The spectral shift function. The papers of M. G. Kreĭn and their further development. (Russian)* Algebra i Analiz **4** (1992), no. 5, 1–44; translation in St. Petersburg Math. J. **4** (1993), no. 5, 833–870.
- [3] D. Bollé, F. Gesztesy, H. Grosse, W. Schweiger, B. Simon, *Witten index, axial anomaly, and Kreĭn's spectral shift function in supersymmetric quantum mechanics*, J. Math. Phys. **28** (1987), no. 7, 1512–1525.
- [4] C. Callias, *Axial Anomalies and Index Theorems on Open Spaces*, Commun. Math. Phys. **62** (1978), 213–234.
- [5] M. G. Krein, *On the trace formula in perturbation theory (Russian)* Mat. Sb. **33** (75), no. 3 (1953), 597–626.
- [6] L. Martínez Alonso and E. Olmedilla, *Trace identities in the inverse scattering transform method associated with matrix Schrödinger operators*, J. Math. Phys. **23** (1982), no. 11, 2116–2121.
- [7] M. Reed and B. Simon, *Methods of modern mathematical physics, I. Functional analysis*. Academic Press, 1980.
- [8] J. Robbin and D. Salamon, *The spectral flow and the Maslov index*, Bull. London Math. Soc. **27** (1995), no. 1, 1–33.
- [9] D. R. Yafaev, *Mathematical scattering theory. General theory*. American Mathematical Society, Providence, RI, 1992.

DEPARTMENT OF MATHEMATICS, KING'S COLLEGE, LONDON. STRAND, LONDON, WC2R 2LS, U.K.

E-mail address: alexander.pushnitski@kcl.ac.uk